

Super Accelerated Motion in Rindler Spacetime

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Abstract: In the general relativity theory, we discover formulas that the super accelerated matter moves with the acceleration α_0 about Rindler space-time. We can represent the super accelerated motion about coordinates x, ct, ξ^0 . PACS Number:04.04.90.+e

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1. INTRODUCTION

In the general relativity theory, we discover formulas that the super accelerated matter moves with the acceleration α_0 about Rindler space-time.

At first, Rindler coordinate is

$$ct = \left(\frac{c^2}{\alpha_0} + \xi^1 \right) \sinh \left(\frac{\alpha_0 \xi^0}{c} \right) \quad (1)$$

$$x = \left(\frac{c^2}{\alpha_0} + \xi^1 \right) \cosh \left(\frac{\alpha_0 \xi^0}{c} \right) - \frac{c^2}{\alpha_0} \quad (2)$$

$$y = \xi^2, z = \xi^3 \quad (3)$$

In Eq (1),

$$\xi^1 + \frac{c^2}{\alpha_0} = \frac{ct}{\sinh \left(\frac{\alpha_0 \xi^0}{c} \right)} \quad (4)$$

If we insert Eq(4) in Eq(2),

$$x = ct \coth \left(\frac{\alpha_0 \xi^0}{c} \right) - \frac{c^2}{\alpha_0} \quad (5)$$

If we insert Eq(5) in Eq(2),

$$ct \coth \left(\frac{\alpha_0 \xi^0}{c} \right) = \left(\frac{c^2}{\alpha_0} + \xi^1 \right) \cosh \left(\frac{\alpha_0 \xi^0}{c} \right) \quad (6)$$

Hence, the result is

$$\xi^1 = \frac{ct}{\sinh \left(\frac{\alpha_0 \xi^0}{c} \right)} - \frac{c^2}{\alpha_0} \quad (7)$$

2. THE SUPER ACCELERATED MOTION ABOUT AN UNIFORMLY ACCELERATED FRAME

$$d\tau^2 = (1 + \frac{a_0 \xi^1}{c^2})^2 (d\xi^0)^2 - \frac{1}{c^2} [(d\xi^1)^2 + (d\xi^2)^2 + (d\xi^3)^2] \quad (8)$$

Hence, if the super accelerated matter moves with the acceleration a_0' about an uniformly accelerated frame,

$$a_0' = \frac{d}{d\xi^0} \left[\frac{\frac{d\xi^1}{d\xi^0}}{\sqrt{(1 + \frac{a_0 \xi^1}{c^2})^2 - (\frac{d\xi^1}{d\xi^0})^2 / c^2}} \right] \quad (9)$$

If we compute,

$$(a_0' \xi^0)^2 = \frac{(\frac{d\xi^1}{d\xi^0})^2}{(1 + \frac{a_0 \xi^1}{c^2})^2 - (\frac{d\xi^1}{d\xi^0})^2 / c^2} \quad (10)$$

Then,

$$(\frac{d\xi^1}{d\xi^0})^2 = (a_0' \xi^0)^2 [(1 + \frac{a_0 \xi^1}{c^2})^2 - (\frac{d\xi^1}{d\xi^0})^2 / c^2] \quad (11)$$

If we compute about $\frac{d\xi^1}{d\xi^0}$,

$$\frac{d\xi^1}{d\xi^0} = \frac{a_0' \xi^0 (1 + \frac{a_0 \xi^1}{c^2})}{\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}}} \quad (12)$$

In Eq(12), if we multiply $\frac{d\xi^0}{(1 + \frac{a_0 \xi^1}{c^2})}$,

$$\frac{d\xi^1}{1 + \frac{a_0 \xi^1}{c^2}} = \frac{a_0' \xi^0 d\xi^0}{\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}}} \quad (13)$$

If we integrate Eq(13),

$$\frac{c^2}{a_0} \ln \left| 1 + \frac{a_0 \xi^1}{c^2} \right| = \frac{c^2}{a_0} \left(\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}} - 1 \right) \quad (14)$$

Hence, if we compute about the coordinate ξ^1 , we can represent the super accelerated motion by

Rindler coordinates ξ^1, ξ^0 .

$$\xi^1 = \frac{c^2}{a_0} \left[\exp \frac{a_0}{a_0} \left(\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}} - 1 \right) - 1 \right] \quad (15)$$

If we insert Eq(7) in Eq(15),

$$\xi^1 = \frac{c^2}{a_0} \left[\exp \frac{a_0}{a_0} \sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}} - 1 \right] = \frac{ct}{\sinh(\frac{a_0 \xi^0}{c})} - \frac{c^2}{a_0} \quad (16)$$

Hence, we can represent the super accelerated motion about coordinates ct, ξ^0 .

$$ct = \frac{c^2}{a_0} \exp \frac{a_0}{a_0} \left(\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}} - 1 \right) \sinh \left(\frac{a_0 \xi^0}{c} \right) \quad (17)$$

If we insert Eq(17) in Eq(5), we can represent the super accelerated motion about coordinates x, ξ^0 .

$$x = ct \coth \left(\frac{a_0 \xi^0}{c} \right) - \frac{c^2}{a_0} = \frac{c^2}{a_0} \exp \frac{a_0}{a_0} \left(\sqrt{1 + \frac{(a_0' \xi^0)^2}{c^2}} - 1 \right) \cosh \left(\frac{a_0 \xi^0}{c} \right) - \frac{c^2}{a_0} \quad (18)$$

3. CONCLUSION

In the general relativity theory, we discover formulas that the super accelerated matter moves with the acceleration in an uniformly accelerated frame.

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